

# A highly sensitive chemical gas detecting transistor based on highly crystalline CVD-grown MoSe<sub>2</sub> films

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Received: 21 July 2016

Revised: 24 August 2016

Accepted: 14 September 2016

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## KEYWORDS

transition metal dichalcogenides, MoSe<sub>2</sub>, chemical sensors, chemical vapor deposition

## ABSTRACT

Layered semiconductors with atomic thicknesses are becoming increasingly important as active elements in high-performance electronic devices owing to their high carrier mobilities, large surface-to-volume ratios, and rapid electrical responses to their surrounding environments. Here, we report the first implementation of a highly sensitive chemical-vapor-deposition-grown multilayer MoSe<sub>2</sub> field-effect transistor (FET) in a NO<sub>2</sub> gas sensor. This sensor exhibited ultra-high sensitivity ( $S = \text{ca. } 1,907$  for NO<sub>2</sub> at 300 ppm), real-time response, and rapid on-off switching. The high sensitivity of our MoSe<sub>2</sub> gas sensor is attributed to changes in the gap states near the valence band induced by the NO<sub>2</sub> gas absorbed in the MoSe<sub>2</sub>, which leads to a significant increase in hole current in the off-state regime. Device modeling and quantum transport simulations revealed that the variation of gap states with NO<sub>2</sub> concentration is the key mechanism in a MoSe<sub>2</sub> FET-based NO<sub>2</sub> gas sensor. This comprehensive study, which addresses material growth, device fabrication, characterization, and device simulations, not only indicates the utility of MoSe<sub>2</sub> FETs for high-performance chemical sensors, but also establishes a fundamental understanding of how surface chemistry influences carrier transport in layered semiconductor devices.

## 1 Introduction

In recent years, the sensing of biological and chemical

molecules has become increasingly important for environmental, biomedical, and agricultural applications [1–4]. NO<sub>2</sub> is one of the most dangerous environ-

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mentally hazardous gas species that we encounter in our daily lives [5]. It has been reported that long-term exposure to NO<sub>2</sub> gas can lead to lung diseases and an increased risk of respiratory symptoms, with the exposure time to NO<sub>2</sub> gas for humans being limited to less than 8 h, even at 3 ppm [6]. In addition, the NO<sub>2</sub> gas generated from automobiles and power plants is thought to contribute to nitric acid rain, which accelerates the degradation of the environment. Consequently, there is an increasing demand for highly sensitive and reliable NO<sub>2</sub> gas sensors to protect human health and to provide timely alarms in the chemical industry.

Field-effect transistors (FETs), which are basic building blocks in electronic circuits, have recently emerged as a novel sensing technology [5, 7–10]. Although there are many resistor-type sensors that can detect various harmful gases [11–13], FET-based gas sensors are generally superior to the simple resistor-type sensors owing to their high sensitivity, large signal-to-noise ratio, short response time, and reliability. The high-sensitivity response of FET sensors is conferred by the gate electrode, which provides extra freedom for current modulation. The sensitivity of FETs can also be evaluated from other parameters, such as subthreshold swing, threshold voltage, and variation of current.

A wide variety of semiconductors have been investigated as active materials for NO<sub>2</sub> sensors, including amorphous organic semiconductors, porous silicon, silicon nanowires, carbon nanotubes, and metal oxide nanowires [5, 14–18]. However, layered semiconductors such as transition metal dichalcogenides (TMDs) have been very rarely investigated for NO<sub>2</sub> sensing, especially in the form of FET-type sensors, despite two-dimensional (2D) materials being promising candidates for chemical sensors owing to their large surface-to-volume ratios [8, 11–13, 19–21].

In the present work, we report the first example of a highly sensitive FET-type NO<sub>2</sub> gas sensor based on a highly crystalline MoSe<sub>2</sub> film prepared by chemical vapor deposition (CVD). When the device is exposed to different NO<sub>2</sub> environments, the off-current of the MoSe<sub>2</sub> transistor increases exponentially with NO<sub>2</sub> gas concentration, exhibiting a maximum sensitivity (*S*) of ca. 1,907 at 300 ppm. Device modeling and quantum

transport simulations along with the measurement data indicate that the high sensitivity of MoSe<sub>2</sub> FETs to NO<sub>2</sub> gas may be attributed to the large variation in the hole current caused by changes in the gap states in the MoSe<sub>2</sub>, which are, in turn, induced by the absorbed NO<sub>2</sub> gas. Our comparative analysis of FET-based and diode-based sensors demonstrates the considerable advantages of using a gate electrode to boost the sensitivity of devices. Finally, we demonstrate the real-time response and fast switching behavior of a FET-based device, highlighting the robustness and potential of our MoSe<sub>2</sub> FET for NO<sub>2</sub> gas sensors.

## 2 Experimental

### 2.1 Synthesis of bulk MoSe<sub>2</sub>

Bulk MoSe<sub>2</sub> was grown by a modified CVD method [22]. First, Mo powder (99.95% purity) and Se shot (99.999% purity) were used under low pressure (10<sup>-5</sup> Torr) to form polycrystalline MoSe<sub>2</sub> compounds. Then, the polycrystalline MoSe<sub>2</sub> and Mo-coated Si substrate precursors were heated to 1,050 and 1,000 °C, respectively, to synthesize bulk MoSe<sub>2</sub>. The MoSe<sub>2</sub> grains exhibited hexagonal structures with sizes in the micrometer range.

### 2.2 Material characterization

The CVD-grown multilayer MoSe<sub>2</sub> samples were characterized using atomic force microscopy (AFM), X-ray diffraction (XRD), Raman spectroscopy, and transmission electron microscopy (TEM). The topography of MoSe<sub>2</sub> flakes deposited on SiO<sub>2</sub>/Si substrates was measured using a PSIA XE7 atomic force microscope in non-contact mode. XRD patterns were recorded on a Bruker D8 Advance X-ray diffractometer using Cu K $\alpha$  radiation. Raman spectra were measured using a high-resolution Raman spectrometer (Invid Raman microscope, Renishaw) with laser excitation at 514.5 nm. The TEM images were obtained using a Tecnai F20 FEI transmission electron microscope at an accelerating voltage of 200 kV.

### 2.3 Device fabrication and measurements

As-grown MoSe<sub>2</sub> was mechanically exfoliated and transferred to an Si/SiO<sub>2</sub> surface. Then, the Si wafer

was immersed in acetone and isopropyl alcohol for 2 h to remove residues. The device was fabricated using conventional photolithography and lift-off methodology. The Ti/Au (20 nm/300 nm) bilayer metal of the source and drain electrodes was deposited using an e-beam evaporator. After patterning, the device was annealed at 200 °C for 2 h in an H<sub>2</sub> environment. The electrical properties of the device were measured using a Keithley 4200 semiconductor characterization system.

## 2.4 Sensor measurement

The fabricated device was wire-bonded to a leadless chip carrier to probe the sample in a vacuum chamber. The schematic of the experimental setup is shown in Fig. S1 (in the Electronic Supplementary Material (ESM)). The sensing experiments were performed by exposing the MoSe<sub>2</sub> device to NO<sub>2</sub> diluted with N<sub>2</sub>. The NO<sub>2</sub> concentration of the mixed gas was controlled by adjusting the flow rates of the N<sub>2</sub> and NO<sub>2</sub> by means of mass flow controllers. All the gas-sensing experiments and characterizations were performed at room temperature.

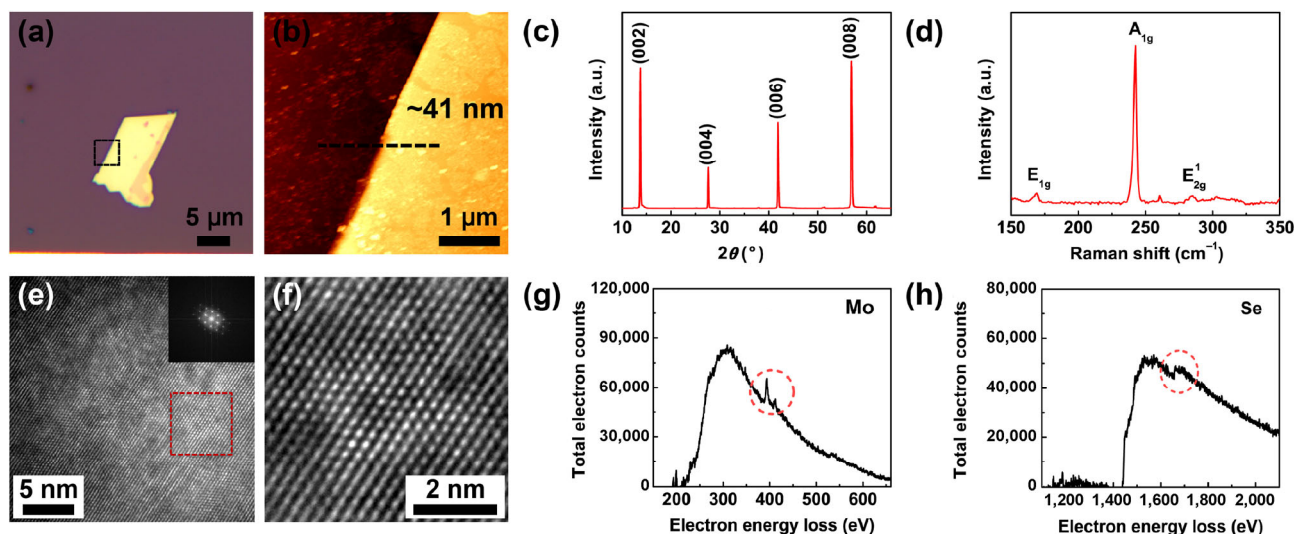
## 2.5 Device modeling and simulations

The transport characteristics were simulated using the non-equilibrium Green's function (NEGF) formalism with an effective mass approximation, self-consistently with Poisson's equation [23]. Since the device shows ambipolar behavior, the electron (n-type) current ( $I_{ds,n}$ ) and hole (p-type) current ( $I_{ds,p}$ ) were calculated separately, and the total current was obtained as  $I_{ds} = I_{ds,n} + I_{ds,p}$ . Values of  $0.55m_0$  and  $0.71m_0$  were used for electron and hole effective mass, respectively [24], where  $m_0$  is free electron mass. The following parameters were used for the nominal device: a 15 nm intrinsic monolayer MoSe<sub>2</sub> channel ( $\epsilon_r = 6.9$ ) [25], and a single-gate device structure with a 2.6 nm-thick SiO<sub>2</sub> gate dielectric. Ballistic transport was assumed. These simulation settings can be used to explore the underlying physics, and will not alter our conclusions for larger devices. Schottky barrier (SB) heights for electrons and holes are chosen such that  $\Phi_{Bn} + \Phi_{Bp} = E_g$  for multilayer MoSe<sub>2</sub> (0.84 eV) [26] in the absence of NO<sub>2</sub> gas. As the NO<sub>2</sub> gas concentration was increased, the effective  $\Phi_{Bp}$  was varied from 0.74 to 0.34 eV in

order to address the extra leakage paths for holes without changing  $\Phi_{Bn}$  (0.1 eV), and channel doping was also increased (up to  $8 \times 10^{18} \text{ cm}^{-3}$  for a  $\Phi_{Bp}$  of 0.34 eV) to take the p-doping effect into account.

## 3 Results and discussion

Multilayer MoSe<sub>2</sub> was prepared by CVD onto a SiO<sub>2</sub> substrate. The details of the synthesis were reported in our previous work [22]. Figures 1(a) and 1(b) show optical microscopy and AFM images of a CVD-grown MoSe<sub>2</sub> flake. The thickness of the MoSe<sub>2</sub> flake is ca. 41 nm. Since the quality and crystallinity of materials directly affect device performance, we characterized the multilayer MoSe<sub>2</sub> using XRD, Raman spectroscopy, and TEM. In the XRD pattern (Fig. 1(c)), prominent peaks appear at  $2\theta$  values of 13.72°, 27.65°, 41.96°, and 57.01°, which correspond to the (002), (004), (006), and (008) planes, respectively. This reveals that the MoSe<sub>2</sub> has a hexagonal, single-crystalline structure [27]. The strongest peak at  $2\theta = 13.72^\circ$ , corresponding to the (002) plane, indicates the preferential growth of the MoSe<sub>2</sub> crystallites in this direction. The Raman spectrum of the MoSe<sub>2</sub> (Fig. 1(d)) presents several typical peaks. An intense out-of-plane A<sub>1g</sub> peak appears at 242 cm<sup>-1</sup>, and weak in-plane E<sub>1g</sub> and E<sub>2g</sub> peaks appear at 169 and 285 cm<sup>-1</sup>, respectively [28, 29]. A plan-view high-resolution TEM image of the as-grown multilayer MoSe<sub>2</sub> is shown in Fig. 1(e). The inset in Fig. 1(e) is a fast Fourier transform (FFT) pattern taken from the whole image. Figure 1(f) is an inverse FFT from the red dotted square in Fig. 1(e) obtained from the primary reflections only. Figures 1(e) and 1(f) clearly demonstrate that the as-synthesized film has a highly crystalline structure and hexagonal symmetry [30, 31]. This is in good agreement with the bulk characterization results obtained by XRD. In the electron energy loss spectra shown in Figs. 1(g) and 1(h), the Mo M<sub>4,5</sub> edge at 227 eV (the M<sub>3</sub> and M<sub>2</sub> edges are at 392 and 410 eV, respectively) and the Se L<sub>3</sub> and L<sub>2</sub> edges at 1,436 and 1,476 eV, respectively (the L<sub>1</sub> edge is at 1,654 eV) can be clearly identified, indicating that the as-synthesized film contains Mo and Se. It should be noted that the O K edge at 532 eV, shown in Fig. 1(g), is negligible. This confirms that the MoO<sub>3</sub> precursor is completely



**Figure 1** (a) Optical image, (b) AFM image, (c) XRD pattern, and (d) Raman spectrum of a CVD-grown multilayer MoSe<sub>2</sub> flake. The image in (b) was taken from the dashed black square in (a), and the thickness of the MoSe<sub>2</sub> flake is about 41 nm. (e) Plan-view high-resolution TEM image of the as-synthesized MoSe<sub>2</sub> film. Inset in (e) is an FFT from the main panel. (f) Inverse FFT image obtained from the primary reflections in the FFT from the dashed red square in (e). Background-subtracted electron energy loss spectra showing (g) the Mo M<sub>4,5</sub> edge at 227 eV (Mo M<sub>3</sub> and M<sub>2</sub> edges at 392 and 410 eV loss are shown in the red dotted circle in (g)) and (h) the Se L<sub>3</sub> and L<sub>2</sub> edges at 1,436 and 1,476 eV (Se L<sub>1</sub> edge at 1654 eV loss is shown in the red dotted circle in (h)).

transformed to MoSe<sub>2</sub> during the CVD process.

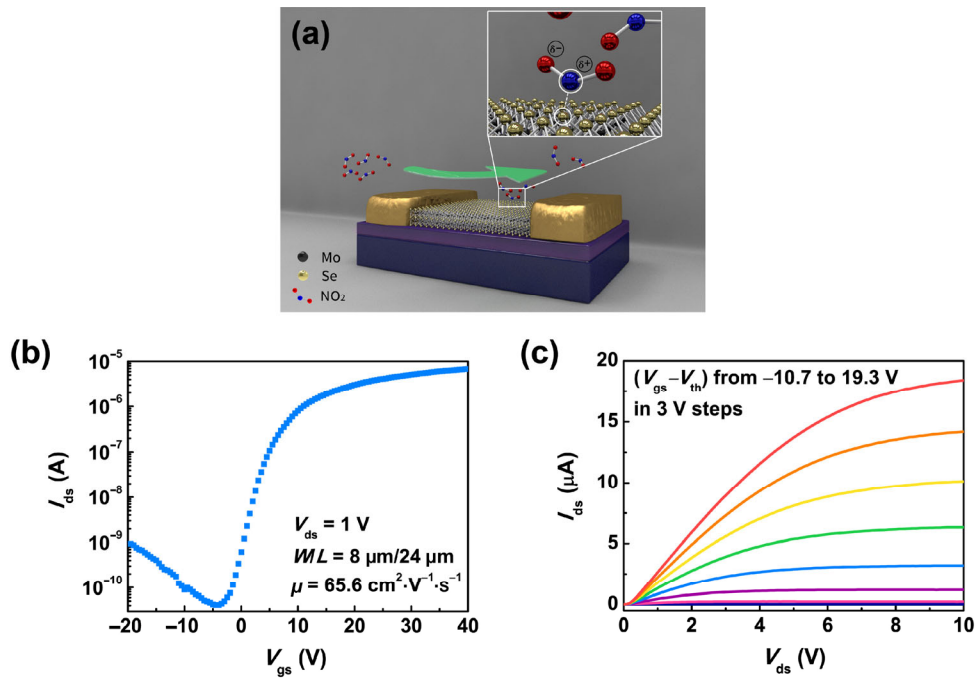
The CVD-grown multilayer MoSe<sub>2</sub> samples were used as an active channel in back-gated FETs to detect NO<sub>2</sub> gas. The as-grown bulk MoSe<sub>2</sub> was mechanically exfoliated to obtain multilayer MoSe<sub>2</sub> flakes, and then Ti/Au (20 nm/300 nm) was deposited onto it to prepare source and drain electrodes. An optical image of the fabricated MoSe<sub>2</sub> FET is shown in Fig. S2(a) (in the ESM). The thickness of the MoSe<sub>2</sub> flake is about 65 nm, as measured using AFM (Figs. S2(b) and S2(c) in the ESM). For gas-sensing measurements, the electrodes of the device were wire-bonded to gold pads on a leadless chip carrier using gold nanowire, and the device was installed in a gas-sensing chamber. Figure 2(a) shows a three-dimensional (3D) schematic diagram of a MoSe<sub>2</sub> transistor-based NO<sub>2</sub> gas sensor. First, we measured the electrical characteristics of the MoSe<sub>2</sub> FET. Typical transfer characteristics ( $I_{ds}-V_{gs}$ ) of the device at  $V_{ds} = 1.0$  V are shown in Fig. 2(b). The device exhibits n-type-dominant ambipolar behavior, indicating that the Ti/Au electrodes form a Schottky barrier such that the barrier height for electrons,  $\Phi_{Bn}$ , is smaller than that for holes ( $\Phi_{Bp}$ ), with an on/off current ratio ( $I_{on}/I_{off}$ ) > 10<sup>5</sup> and a threshold voltage ( $V_{th}$ ) of 6.7 V.

The output characteristics ( $I_{ds}-V_{ds}$ ) of the same device are shown in Fig. 2(c). The  $I_{ds}-V_{ds}$  curves show linearity at low drain voltages, indicating the formation of an Ohmic contact at room temperature, and exhibit clear saturation behavior at high drain voltages. The field-effect mobility ( $\mu$ ) is estimated to be 65.6 cm<sup>2</sup>·V<sup>-1</sup>·s<sup>-1</sup>, which is calculated from the triode region using Eq. (1):

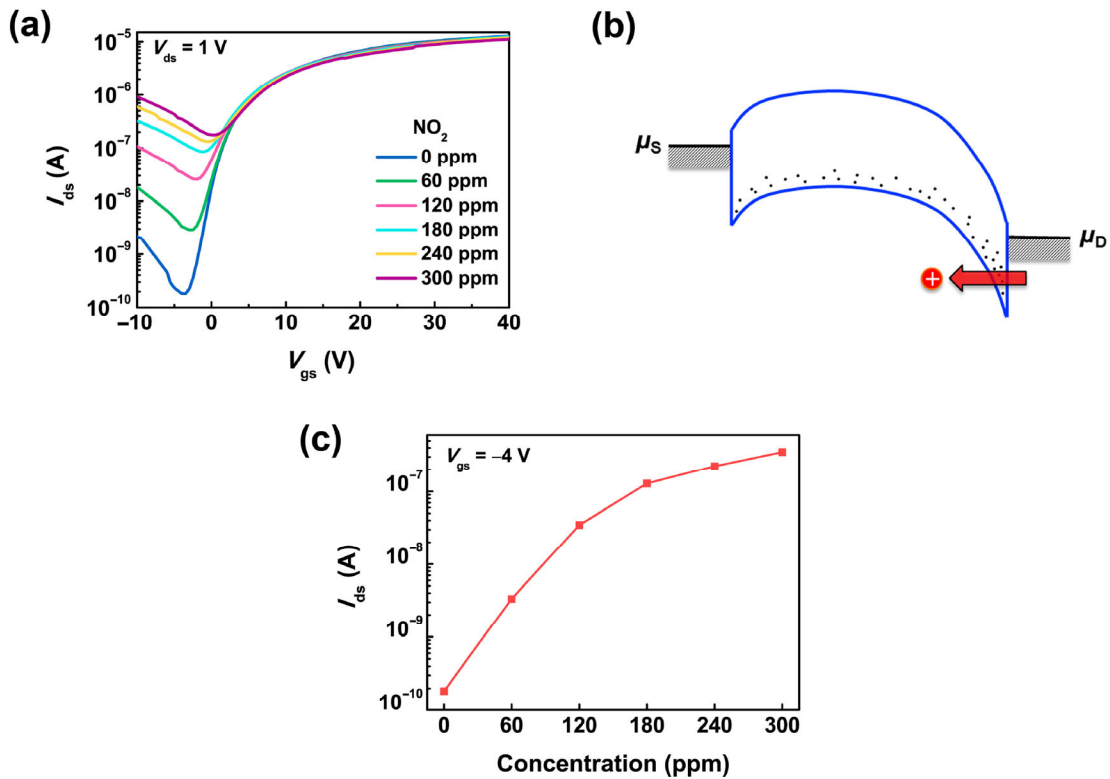
$$\mu = \left( \frac{L}{WC_{ox}V_{ds}} \right) \left( \frac{\partial I_{ds}}{\partial V_{gs}} \right) \quad (1)$$

where  $L$  is the device channel length (24 μm),  $W$  is the device channel width (8 μm), and  $C_{ox}$  is the gate capacitance per unit area ( $1.15 \times 10^{-8}$  F·cm<sup>-2</sup>).

It was expected that a MoSe<sub>2</sub> transistor-based gas sensor would exhibit an ultra-high sensing capacity due to the large surface-to-volume ratio of 2D materials. Figure 3(a) shows the transfer characteristics ( $I_{ds}-V_{gs}$ ) at  $V_{ds} = 1$  V under different NO<sub>2</sub> gas concentrations. The curve measured in an N<sub>2</sub> environment in the absence of NO<sub>2</sub> exhibits similar characteristics to that in air (Fig. 2(b)). When the device is exposed to NO<sub>2</sub> gas, the off-current in the p-branch (i.e., the hole current) increases significantly with NO<sub>2</sub> gas concentration,



**Figure 2** (a) 3D schematic of a MoSe<sub>2</sub> FET-based NO<sub>2</sub> gas sensor. (b) Transfer characteristic ( $I_{ds}$ – $V_{gs}$ ) of a multilayer MoSe<sub>2</sub> transistor at  $V_{ds} = 1$  V in a logarithmic scale. (c) Output characteristics ( $I_{ds}$ – $V_{ds}$ ) at different ( $V_{gs}$ – $V_{th}$ ) from  $-10.7$  to  $19.3$  V in  $3$  V steps.

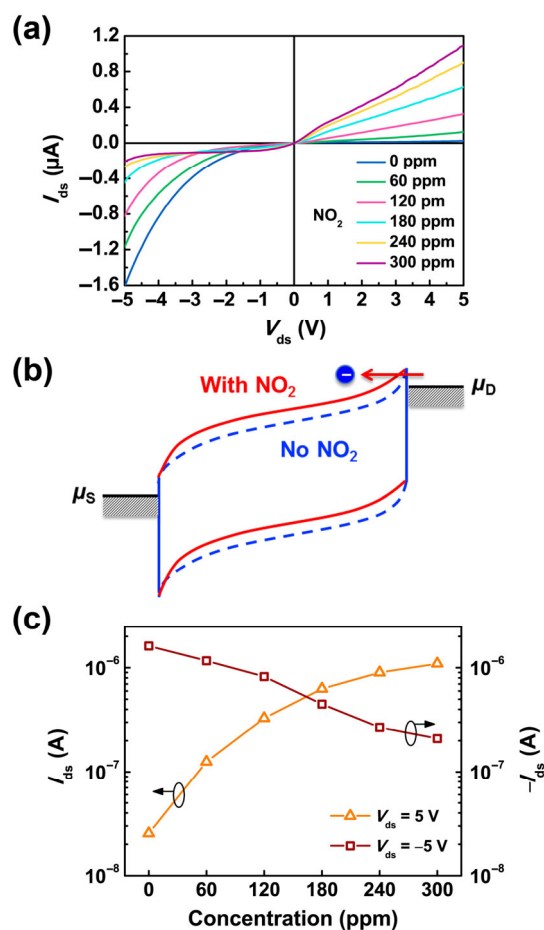


**Figure 3** (a) Transfer characteristics ( $I_{ds}$ – $V_{gs}$ ) measured at NO<sub>2</sub> concentrations of 0, 60, 120, 180, 240, and 300 ppm plotted on a logarithmic scale. (b) Energy band diagram of MoSe<sub>2</sub> FET under NO<sub>2</sub> gas at  $V_{ds} > 0$  and  $V_{gs} < 0$ .  $\mu_{S,D}$  are the chemical potentials in the source and drain. (c)  $I_{ds}$  vs. NO<sub>2</sub> concentration at  $V_{gs} = -4$  V, extracted from (a).

whereas the electron current in the n-branch is largely constant (Fig. 3(a)). It has been reported that  $\text{NO}_2$  absorbed in TMDs such as  $\text{MoSe}_2$ ,  $\text{MoS}_2$ , and  $\text{WSe}_2$  leads to states in the band gap region close to the valence band [32, 33], and the same was expected for our device.

These states induced by  $\text{NO}_2$  molecules can provide extra tunneling paths for hole conduction, as depicted in Fig. 3(b), leading to the hole current increasing as the device is exposed to  $\text{NO}_2$  gas. This transport mechanism can also be understood and modeled with the suppression of the effective SB for holes with  $\text{NO}_2$ , which will be discussed in detail later. Figure 3(c) plots the  $I_{\text{ds}}$  as a function of  $\text{NO}_2$  gas concentration at  $V_{\text{gs}} = -4$  V. As the  $\text{NO}_2$  concentration increases from 0 to 300 ppm, the current is enhanced by a factor of 1,909. This is the maximum current enhancement observed in the off-state regime, and demonstrates the superior sensing capacity of our  $\text{MoSe}_2$  FETs.

In order to verify the use of FET-based  $\text{MoSe}_2$  sensors for ultra-high sensitivity gas detection, we have also measured the  $I_{\text{ds}}-V_{\text{ds}}$  characteristics of a diode-like device without applying  $V_{\text{gs}}$  for different  $\text{NO}_2$  gas concentrations, as shown in Fig. 4(a). As the concentration of  $\text{NO}_2$  gas increases, the current ( $I_{\text{ds}}$ ) increases in the positive voltage regions but decreases in the negative voltage regions. When a positive voltage is applied on the drain side ( $V_{\text{ds}} > 0$ ), the band diagram of the device is analogous to that shown in Fig. 3(b), where hole current increases significantly with  $\text{NO}_2$  gas owing to the additional band gap states located near to the valence band, as explained above. Conversely, if the opposite voltage is applied across the device ( $V_{\text{ds}} < 0$ ), the chemical potential in the drain becomes higher than that in the source, as shown in Fig. 4(b), where the majority of current is due to electron flow from the drain to the source. In the presence of  $\text{NO}_2$  gas, which can enhance charge transfer as a p-dopant, the conduction and the valence band of  $\text{MoSe}_2$  are shifted upward, causing a wider SB and, thereby, reducing the electron current significantly [8, 12, 32]. Although the same trend can be observed for the on-current (i.e., the electron current in the n-branch) in Fig. 3(a), the impact of charge transfer or doping is minimal when a positive



**Figure 4** (a)  $I_{\text{ds}}-V_{\text{ds}}$  for  $\text{NO}_2$  concentrations from 0 to 300 ppm for a diode-like device with no  $V_{\text{gs}}$  applied. (b) Band diagram of the diode-like  $\text{MoSe}_2$  device at  $V_{\text{ds}} < 0$ . Solid red lines show the effect of p-doping with  $\text{NO}_2$  gas. (c)  $I_{\text{ds}}$  vs.  $\text{NO}_2$  concentration at  $V_{\text{ds}} = -5$  and 5 V.

$V_{\text{ds}}$  is applied. Figure 4(c) shows the current as a function of the  $\text{NO}_2$  gas concentration, measured at  $V_{\text{ds}} = -5$  and 5 V for the same device as that in Fig. 4(a). For a change in the  $\text{NO}_2$  concentration from 0 to 300 ppm, the electron current decreases by 87% at  $V_{\text{ds}} = -5$  V and the hole current increases by a factor of 43 at  $V_{\text{ds}} = 5$  V.

Notably, these values are significantly lower than those for a change of current in FET-based  $\text{MoSe}_2$  sensors, where the current is enhanced by a factor of 1,909 for the same variation in  $\text{NO}_2$  concentration, as we have shown above. In addition, Figs. 4(a) and 4(c) clearly indicate that the hole current at positive  $V_{\text{ds}}$  is a more effective indicator for  $\text{NO}_2$  gas using  $\text{MoSe}_2$  sensors.

For a more in-depth understanding of the characteristics of MoSe<sub>2</sub> FET-based sensors, we performed quantum transport simulations using the NEGF method. We modeled the absorption of NO<sub>2</sub> gas in MoSe<sub>2</sub> FETs considering (i) the suppression of effective SB height for holes ( $\Phi_{Bp}$ ) due to the extra states and the increased tunneling conferred by the absorbed NO<sub>2</sub>, and (ii) the effect of p-doping in the channel region. The effective  $\Phi_{Bp}$  was varied from 0.74 to 0.34 eV to minimize the hole-leakage paths for NO<sub>2</sub> gas, without changing the SB height for electrons ( $\Phi_{Bn} = 0.1$  eV). Figure 5(a) shows the simulated  $I_{ds}$ - $V_{gs}$  characteristics in the absence of NO<sub>2</sub> gas and under low and high concentrations of NO<sub>2</sub> gas. They exhibit the same trend as that presented by the experimental data (Fig. 3(a)). The ratio of  $I_{gas}$  to  $I_0$  as a function of  $\Phi_{Bp}$  is plotted in Fig. 5(b) (shown in star symbols), where  $I_{gas}$  is the current at different NO<sub>2</sub> concentrations and  $I_0$  is the current without NO<sub>2</sub> gas, obtained from the simulations at  $V_{gs} = -0.3$  V, which is the voltage resulting in the minimum leakage current for  $I_0$ .  $I_{gas}/I_0$  exhibits a clear exponential increase with decreasing  $\Phi_{Bp}$ , according to the following relationship

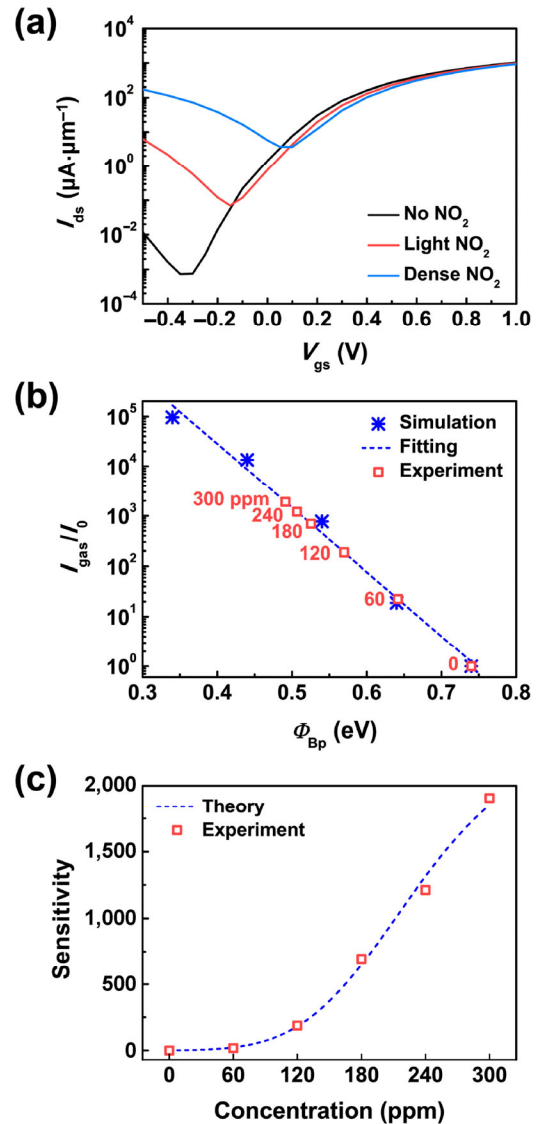
$$\log_{10} \left( \frac{I_{gas}}{I_0} \right) = p_1 \cdot \Phi_{Bp} + p_2 \quad (2)$$

where  $p_1$  and  $p_2$  are  $-12.8$  and  $9.47$ , respectively. Furthermore, we investigated how the effective  $\Phi_{Bp}$  is affected by NO<sub>2</sub> gas concentration ( $c_{NO_2}$ ) by correlating the theoretical expression (the dashed line in Fig. 5(b)) with the experimental data (the square symbols in Fig. 5(b)), resulting in

$$\Phi_{Bp} = p_3 \cdot c_{NO_2}^3 + p_4 \cdot c_{NO_2}^2 + p_5 \cdot c_{NO_2} + p_6 \quad (3)$$

where  $p_3$ ,  $p_4$ ,  $p_5$ , and  $p_6$  are  $-5.22 \times 10^{-9}$ ,  $5.46 \times 10^{-6}$ ,  $-2.00 \times 10^{-3}$ , and  $0.74$ , respectively. This indicates that the effective  $\Phi_{Bp}$  varies significantly with NO<sub>2</sub> gas concentration, gradually reaching saturation as  $c_{NO_2}$  increases. In addition, a direct relationship between  $\left( \frac{I_{gas}}{I_0} \right)$  and  $c_{NO_2}$  can be obtained from Eqs. (2) and (3) as

$$\frac{I_{gas}}{I_0} = 10^{(p_7 \cdot c_{NO_2}^3 + p_8 \cdot c_{NO_2}^2 + p_9 \cdot c_{NO_2})} \quad (4)$$



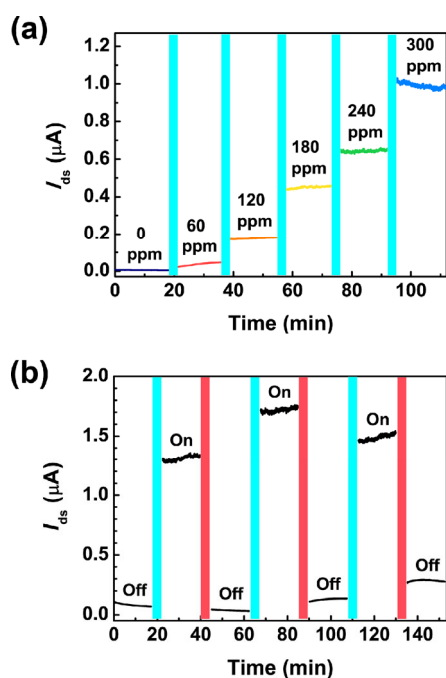
**Figure 5** (a) Simulation results of transfer characteristics under zero, low, and high concentrations of NO<sub>2</sub> gas. (b)  $I_{gas}/I_0$  vs. effective SB height for holes ( $\Phi_{Bp}$ ). The blue stars represent simulation results with various  $\Phi_{Bp}$  obtained at  $V_{gs} = -0.3$  V. The dashed blue line is a linear fit to the simulation results. The red squares represent the experimental data with different gas concentrations, taken from Fig. 3(a) at  $V_{gs} = -4$  V. (c) Sensitivity vs. NO<sub>2</sub> concentration. The dashed blue line is based on theory, and the red square symbols represent the data from the experiments shown in Fig. 3(a) ( $V_{gs} = -4$  V).

where  $p_7$ ,  $p_8$ , and  $p_9$  are  $6.68 \times 10^{-8}$ ,  $-6.99 \times 10^{-5}$ , and  $2.56 \times 10^{-2}$ , respectively.

The sensitivity ( $S$ ) of gas sensors can be defined as  $(I_{gas} - I_0)/I_0$ , which is the most important value in the assessment of gas sensors. We have plotted sensitivity as a function of NO<sub>2</sub> gas concentration in Fig. 5(c)

and Fig. S3 in the ESM. The dashed curve in Fig. 5(c) is a theoretical plot drawn using Eq. (4), and the square symbols are from the experimental data (Fig. 3(a);  $V_{gs} = -4$  V), revealing an excellent match between the theory and experimental observations. Figure 5(c) indicates that our FET-based MoSe<sub>2</sub> sensor effectively detects NO<sub>2</sub> gas, showing a linear response to concentration at >100 ppm, with the maximum sensitivity of ca. 1,907 observed at 300 ppm.

We systematically analyzed the time-resolved response of our MoSe<sub>2</sub> gas sensor. Figure 6(a) shows the real-time current response of a MoSe<sub>2</sub> FET to NO<sub>2</sub> gas, where the concentration of NO<sub>2</sub> gas increases with time. To achieve stable conditions, the different concentrations of NO<sub>2</sub> gas were injected into the chamber for 20 min (light blue regions; not scaled) before the current was measured at  $V_{gs} = -4$  V and  $V_{ds} = 1$  V. The current shows similar behavior to those in



**Figure 6** (a) Real-time current response of MoSe<sub>2</sub> FET to NO<sub>2</sub> gas at different NO<sub>2</sub> concentrations. The light blue regions indicate 20-min NO<sub>2</sub> gas injection (not scaled) into the chamber before the measurement. (b) Real-time on–off switching behavior conducted at a NO<sub>2</sub> concentration of 300 ppm. The light blue and red regions denote injection for 20 min and venting for 30 min (not scaled), respectively. A large negative gate bias ( $V_{gs} = -40$  V) was applied to desorb NO<sub>2</sub> from the surface of the MoSe<sub>2</sub> channel during venting to reduce the recovery time. Current was measured at  $V_{gs} = -4$  V and  $V_{ds} = 1$  V in both (a) and (b).

Fig. 3(a), indicating the robustness of our MoSe<sub>2</sub> gas sensor. As shown in Fig. 6(b), we investigated the on–off switching behavior where the measurement was performed at the same bias conditions as in Fig. 6(a) with an injection gas concentration of 300 ppm. A recovery time of 2 h was typically required under inert N<sub>2</sub> purging gas. However, the recovery time was significantly reduced to 30 min by using a large negative gate bias ( $V_{gs} = -40$  V in Fig. 6(b)). In this case, the repulsive electrical field from the gate bias facilitates the desorption of NO<sub>2</sub> from the MoSe<sub>2</sub> surface, and this observation is consistent with those detailed in previous reports [12, 34]. It is feasible that additional treatments such as UV exposure, thermal annealing, or pulse-gating could be applied to further shorten the recovery time [12, 35–37].

Finally, selectivity is a very important feature of gas sensors. Generally, selectivity can be improved by two different approaches: firstly, the introduction of functional groups [38] or catalytic particles [39, 40], which can react with the target gas; and secondly, the employment of an appropriate filter that can induce catalytic oxidation (i.e., a catalytic filter) of interfering species [41], or physically remove them (i.e., a porous membrane filter) via molecular key-lock systems [42]. Specifically, for NO<sub>2</sub> gas sensors, many different types of membrane-based filters (e.g., carbonaceous chemical filters, carbon nanocones and discs, or nano-carbonaceous filters) have been reported in the literature [43–45]. Therefore, issues of selectivity toward NO<sub>2</sub> gas can be resolved using previously developed filters designed for NO<sub>2</sub> gas sensors placed upstream of the MoSe<sub>2</sub>-based FET sensor.

## 4 Conclusions

We have demonstrated a highly sensitive NO<sub>2</sub> gas sensor based on a CVD-grown MoSe<sub>2</sub> FET, which shows n-type-dominant ambipolar behavior with a field-effect mobility of 65.6 cm<sup>2</sup>·V<sup>-1</sup>·s<sup>-1</sup> and an on/off current ratio of 10<sup>5</sup>. Our experiments showed that the hole current in a MoSe<sub>2</sub> FET significantly increases when it is exposed to NO<sub>2</sub> gas, which is attributed to the extra band gap states induced by the absorbed NO<sub>2</sub> gas. We also performed quantum transport



simulations to model the variation of  $\Phi_{\text{BP}}$  and the effect of p-doping, and the simulations reproduced the experimental observations by considering the induced band gap states. Our study demonstrated the linear response of the sensitivity of the MoSe<sub>2</sub> FET to NO<sub>2</sub> gas concentration in the 100–300 ppm range with a maximum sensitivity of ca. 1,907. In addition, a comparison of FET-based sensors and diode-like MoSe<sub>2</sub> sensors revealed that sensitivity could be significantly enhanced using a negative  $V_{\text{gs}}$ , indicating the great advantage of using FET-based chemical sensors. Finally, we have demonstrated excellent real-time response and on–off switching behavior, verifying the robustness of our device. In this study, our highly sensitive MoSe<sub>2</sub> FET represents a significant advance in 2D layered materials for practical sensing applications.

## Acknowledgements

This research is also supported in part by the National Research Foundation of Korea (Nos. NRF-2013M3C1A3059590, NRF-2014M3A9D7070732, and NRF-2015R1A5A1037548). This work was supported in part by U.S. National Science Foundation under grant CMMI 826276, and in part by NSERC Discovery Grant (No. RGPIN-05920-2014). This research was supported by the Commercializations Promotion Agency for R&D Outcomes (COMPA) funded by the Ministry of Science, ICT and Future Planning (MISP) and in part by the fund from the Korea Institute of Science and Technology (KIST) institutional program. Computing resources were provided by SHARCNET through Compute Canada. D. Y. acknowledges the financial support by WIN Nanofellowship. The authors would like to thank Prof. Jong-Soo Rhyee for supporting CVD MoSe<sub>2</sub> flakes.

**Electronic Supplementary Material:** Supplementary material (schematic of experimental setup, optical and AFM images of MoSe<sub>2</sub> device, and sensitivity vs. NO<sub>2</sub> concentration) is available in the online version of this article at <http://dx.doi.org/10.1007/s12274-016-1291-7>.

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